Edge plasma modeling for divertor configurations with secondary x-points

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Acknowledgments:

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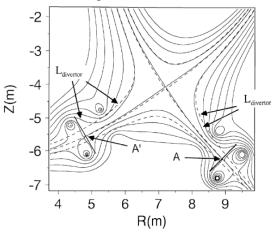


- Divertor configurations with secondary X-points
- Snowflake divertor experiments
- X-Point target divertor
- Topology and grids for edge domain with two x-points
- UEDGE analysis of near-snowflake divertor configurations
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- Summary/Conclusions

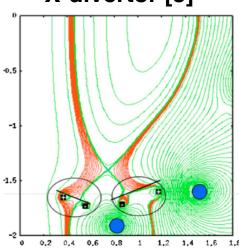
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Configurations with a secondary X-point in divertor considered by many groups in recent years; for example

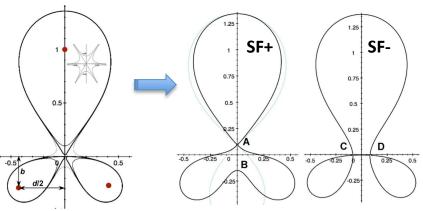
Cusp divertor [1]



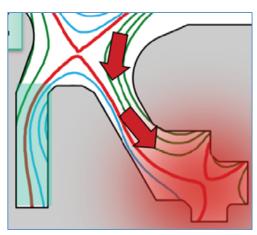
X-divertor [3]



Snowflake divertor [2]



X-point target divertor [4]



[1] H. Takase, J. Phys. Soc. Japan, 70, 609, 2001. [3] M. Kotschenreuther et al., 2004 IAEA FEC, paper IC/P6-43.

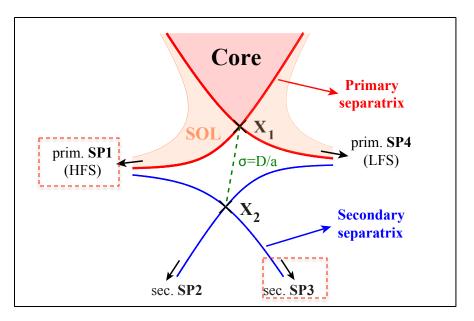
[2] D.D. Ryutov. Phys. Plasmas, 14, 064502, 2007. [4] B. LaBombard et al., Nucl. Fusion 55, 053020, 2015.

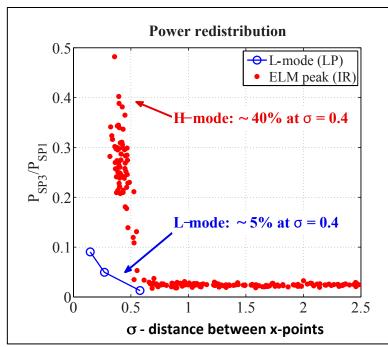


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Experiments on TCV tokamak indicate that enhanced transport zone may exist near the null-point

Figures from Vijvers et al, NF 2014

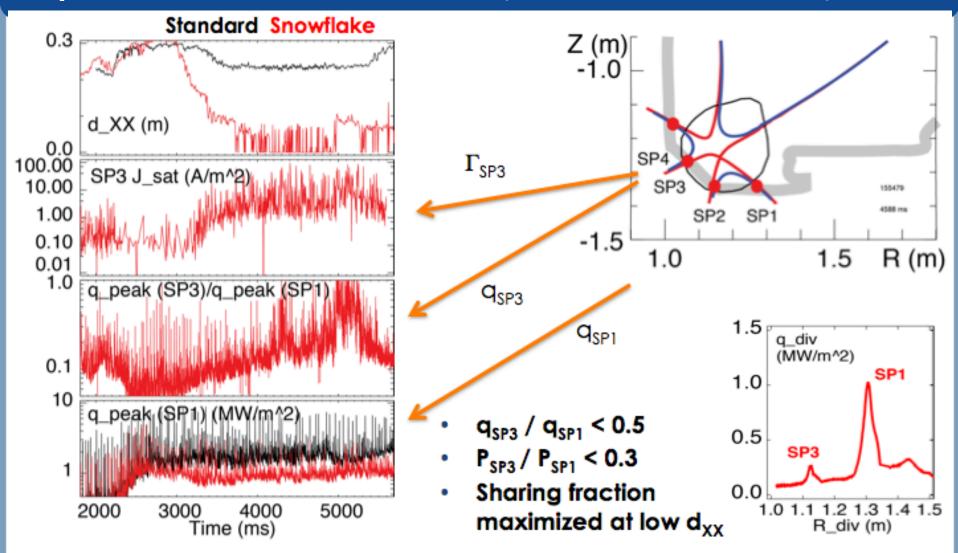




Large fraction of power flows to secondary strike points

- when two X-points get closer
- more during ELM strike

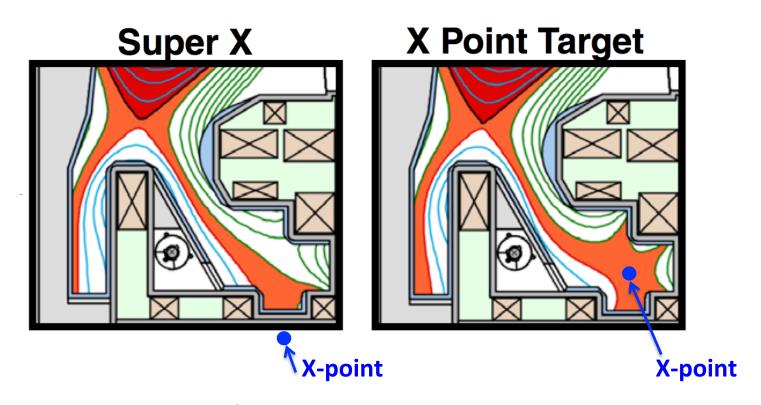
DIII-D: heat and particle fluxes shared among strike points in snowflake divertor (Soukhanovskii – FEC '14)





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X-point target divertor is similar to the super-X divertor, but with the second X-point in the plasma volume



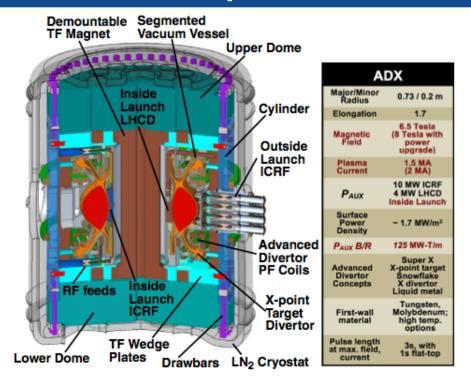
- Like Super-X, exploits 1/R geometric reduction of divertor heat flux
- May produce stable 'X-point MARFE' in the divertor chamber
- Used as a part of the ADX tokamak concept

XPTD: LaBombard et al. 2013 Bull. Am. Phys. Soc. 58 63, and Nucl. Fusion <u>55</u>, 053020, 2015.

SXD: P. Valanju *et al., Phys. Plasmas* <u>16</u>, 056110 (2009)



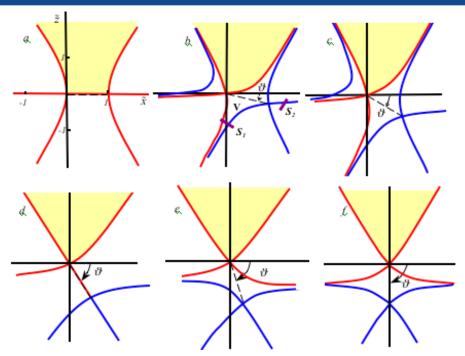
X-point target divertor study is motivated by the ADX tokamak concept discussed at MIT PSFC



- ADX = Advanced Divertor and RF tokamak eXperiment*
- Designed to address critical gaps on pathway to next-step devices
- Advanced divertors
- Advanced RF actuators
- Reactor-prototypical core plasma conditions

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Topological classification of configurations with secondary x-point in the divertor



- Derived from local expansion for inexact snowflake
- Applies to any configuration with secondary x-point
- θ = angle between X-point bisector and horizontal axis
- In addition to shown six cases, there are mirror reflections of cases b,c,d,e

Recent upgrades made in UEDGE include generalization of computational subdomains and mesh generation

θ = angle betweenX-point bisector &horizontal axis

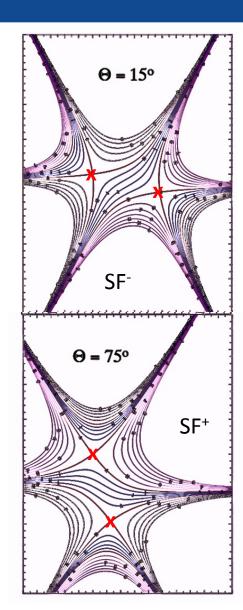
3 mesh regimes

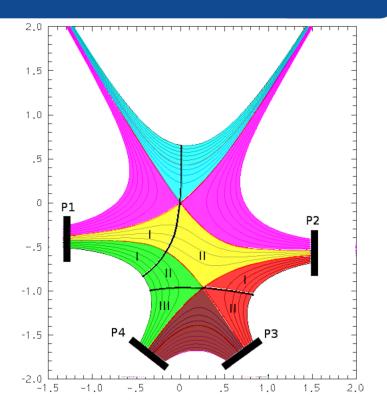
 $0 < \theta < 30^{\circ}$

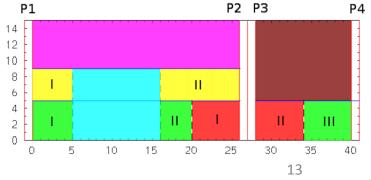
 $30^{\circ} < \theta < 60^{\circ}$

60° < θ < 90°

Unique indexing rules for each regime completed





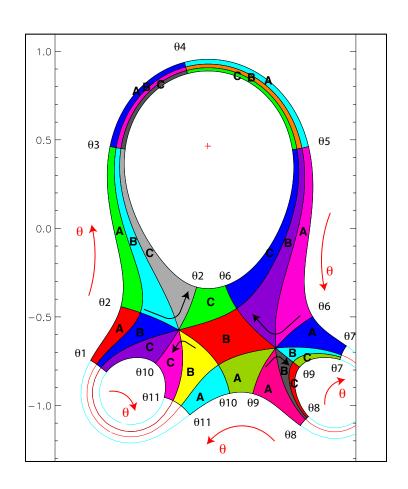


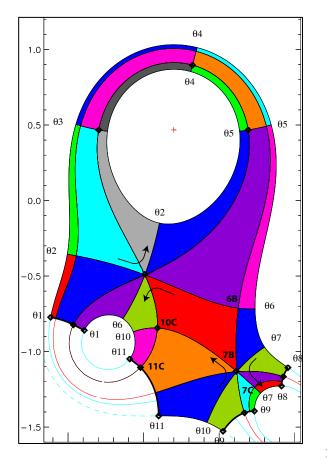


From the point of view of domain topology and grid connectivity there are two distinct SFM regimes

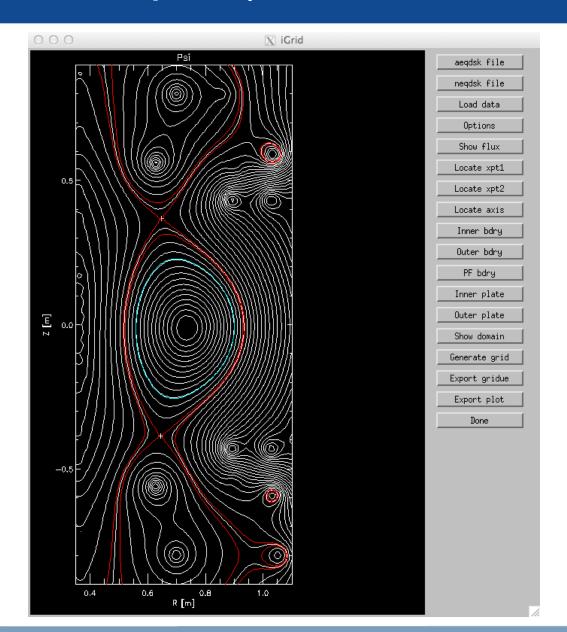
SFM1 – From core boundary there is a path down grad(ψ) to PF boundary

SFM2 – From core boundary down grad(ψ) can only get to SOL boundary





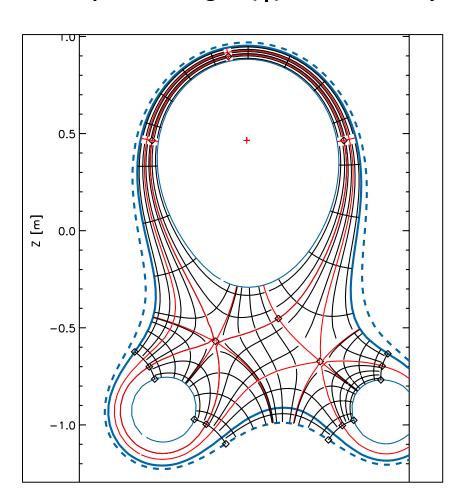
Interactive Grid Generator iGrid (under continuing development) has been used for constructing meshes



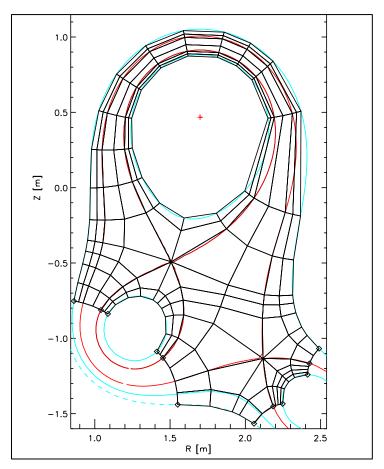
- Variety of flux surface geometry is a challenge for tokamak edge grid generation
- Human eye is still the best tool for recognizing complex patterns
- In iGrid the user guides the code by indicating with the mouse some needed reference points and directions

Orthogonal grids for SFM1 and SFM2 are generated for analytic "3-wire" geometry

SFM1 – From core boundary there is a path down grad(ψ) to PF boundary



SFM2 – From core boundary down grad(ψ) can only get to SOL boundary

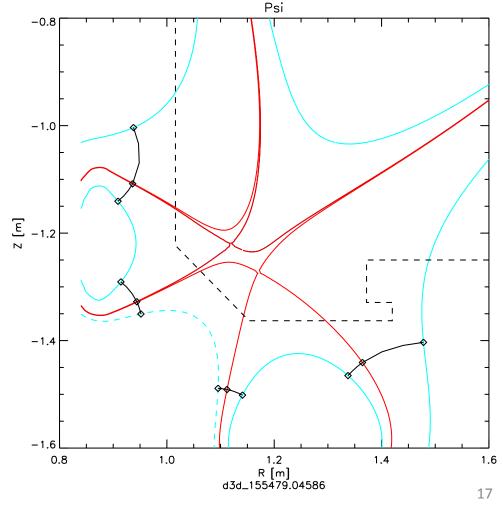


Real tokamak geometry in near-snowflake configuration makes a challenge for grid numerics

Actual DIII-D geometry

1.0 R [m] d3d_155479.04586

Extremely fast convergence of flux surfaces



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UEDGE* (Unified EDGE code) solves a system of fluid equations in axisymmetric tokamak geometry

plasma density

ion II momentum

electron thermal energy

ion thermal energy

neutral density

ad-hoc radial transport

neutral II momentum

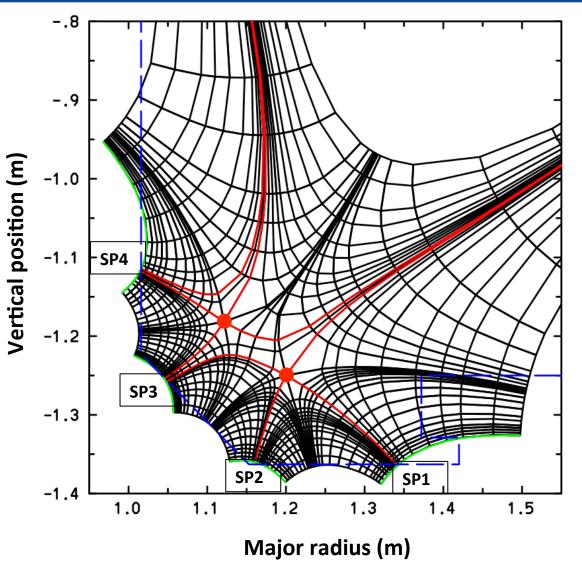
charge conservation

sheath bound. cond.

$$\begin{split} \frac{\partial}{\partial t}(n_{i}) &+ \nabla \bullet (n_{i}\vec{u}_{i}) = -S_{r} + S_{i} \\ n_{i}u_{\perp} &= -D_{\perp i}\nabla_{\perp}n_{i} \\ \frac{\partial}{\partial t}(mn_{i}u_{\parallel i}) &+ \nabla \bullet (mn_{i}u_{\parallel i}\vec{u}_{i} - \eta_{i}\nabla u_{\parallel i}) = -\nabla_{\parallel}P_{i} + mn_{N}n_{i}K_{cx}(u_{\parallel N} - u_{\parallel i}) + mS_{r}u_{\parallel N} - mS_{i}u_{\parallel N} \\ \frac{\partial}{\partial t}(3/2n_{e}T_{e}) &+ \nabla \bullet (\frac{5}{2}n_{e}T_{e}\vec{u}_{e} + \vec{q}_{e}) = \vec{u}_{e} \bullet \nabla (3/2n_{e}T_{e}) - \Pi_{e} \bullet \nabla \vec{u}_{e} + Q_{e} \\ \frac{\partial}{\partial t}(3/2n_{i}T_{i}) &+ \nabla \bullet (\frac{5}{2}n_{i}T_{i}\vec{u}_{i} + \vec{q}_{i}) = \vec{u}_{i} \bullet \nabla (3/2n_{i}T_{i}) - \Pi_{i} \bullet \nabla \vec{u}_{i} + Q_{i} \\ q_{\perp} &= -n\chi_{\perp}\nabla_{\perp}T \\ \frac{\partial}{\partial t}(n_{N}) &+ \nabla \bullet (n_{N}\vec{u}_{N}) = S_{r} - S_{i} \\ \frac{\partial}{\partial t}(mn_{N}u_{\parallel N}) &+ \nabla \bullet (mn_{N}u_{\parallel N}\vec{u}_{N} - \eta_{N}\nabla u_{\parallel N}) = -\nabla_{\parallel}P_{N} - mn_{N}n_{i}K_{cx}(u_{\parallel N} - u_{\parallel i}) - mS_{r}u_{\parallel N} + mS_{i}u_{\parallel N} \\ \frac{\partial}{\partial t}(mn_{N}u_{\parallel N}) &+ \nabla \bullet (mn_{N}u_{\parallel N}\vec{u}_{N} - \eta_{N}\nabla u_{\parallel N}) = -\nabla_{\parallel}P_{N} - mn_{N}n_{i}K_{cx}(u_{\parallel N} - u_{\parallel i}) - mS_{r}u_{\parallel N} + mS_{i}u_{\parallel N} \\ \nabla \bullet J(\phi) &= 0 \\ J_{\parallel} &= \frac{en}{0.51mv} \frac{B_{x}}{B} \left(\frac{1}{n} \frac{\partial P_{e}}{\partial x} - e \frac{\partial \phi}{\partial x} + 0.71 \frac{\partial T_{e}}{\partial x} \right) \\ J_{r} &= \sigma_{\perp} E_{r} \\ \phi &= \frac{-Te}{e} \ln \left[2\sqrt{\pi} \left(\frac{J_{\parallel} - enu_{\parallel i}}{env_{ie}} \right) \right] \end{split}$$

^{*}T. D. Rognlien et al., J. Nucl. Mater. 196–198, 347 (1992)

UEDGE SFM1 mesh for DIII-D shot 155479

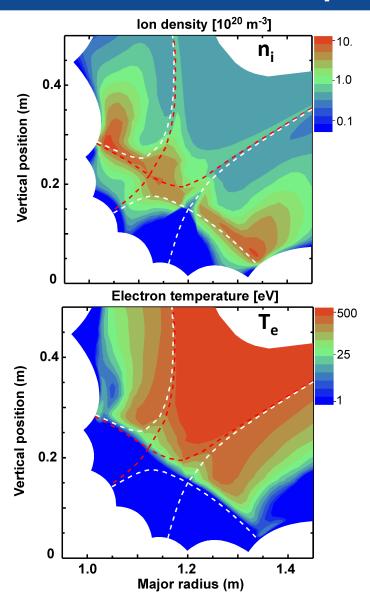


X-point separation/ minor radius

$$\sigma$$
 = 0.1/0.5 = 0.2

50% greater mesh resolution used for simulations

2D UEDGE SF solutions for 2% carbon show strong variations across separatrices; radiation is well spread



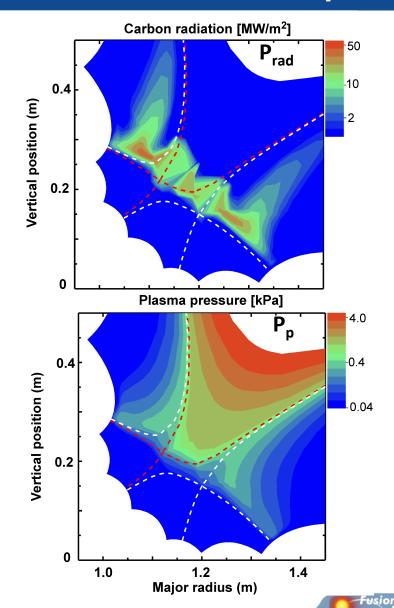
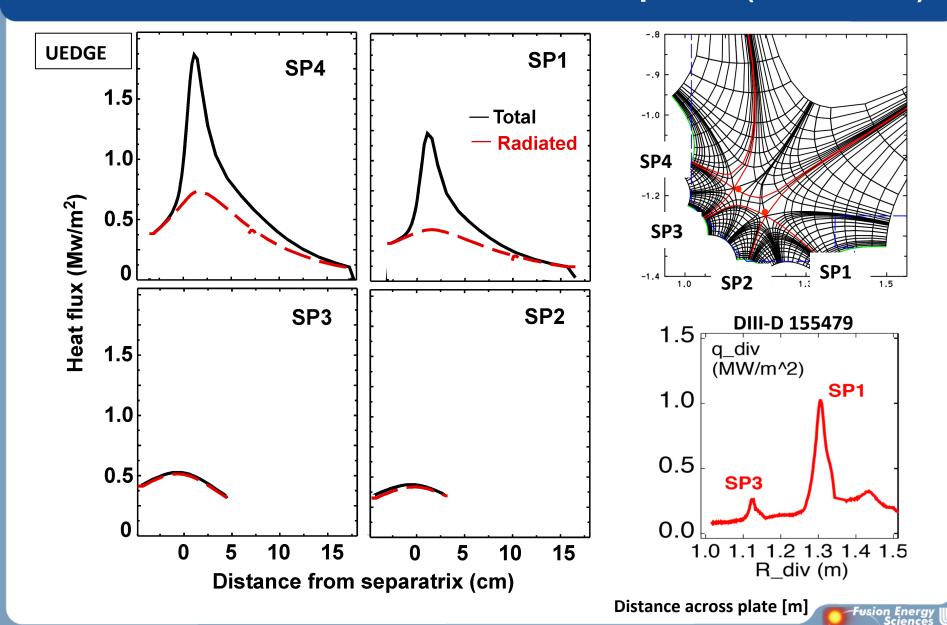


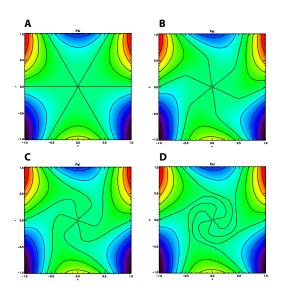
Plate heat-fluxes are < 2 MW/m²; only radiative flux is visible in the middle between two strike points (2% carbon)



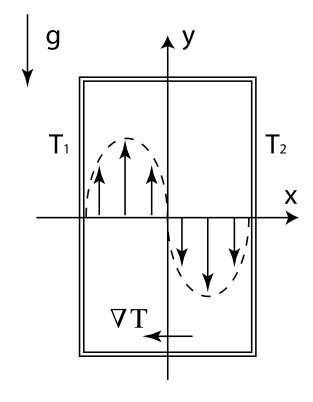
Convective cell formation near null dubbed "the churning mode" may be responsible for heat redistribution*

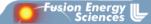
Driven by crossed magnetic curvature and grad(P) near null where poloidal beta is large

Solving plasma fluid equations demonstrates formation of the churning mode**



Similar to thermal convection due to crossed gravity and temperature gradient



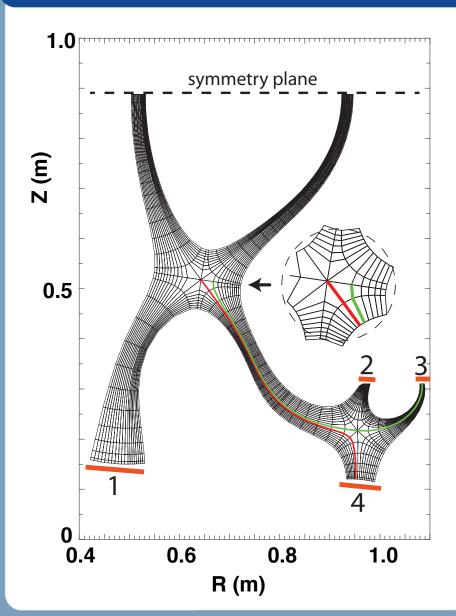


^{*}D.D. Ryutov et al., Physica Scripta 89, 8, 088002 (2014)

^{**} M.V. Umansky and D.D. Ryutov, submitted (2015)

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UEDGE is used to model both X-points in an XPTD for the lower half of up-down symmetric configuration



Mesh constructed in UEDGE by combining two lower-half single-null domains (with scripts developed by M.E. Rensink)

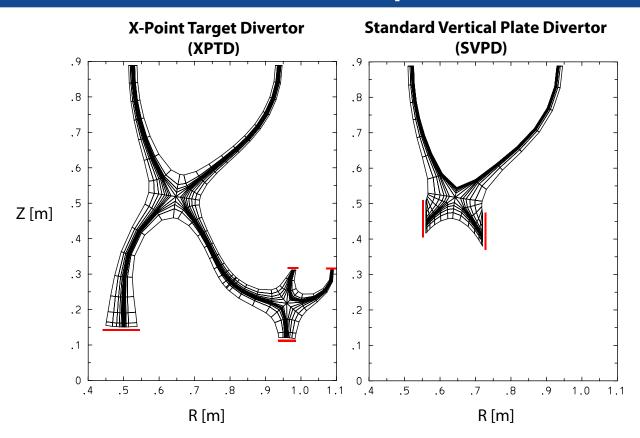
Use UEDGE fluid transport model

- Fluid neutrals (inertial)
- Fixed fraction impurity radiation
- No drifts
- Four orthogonal target plates
- 100% recycling on all walls

Use geometry & parameters from LaBombard et al., NF 2015

- MHD equilibrium provided by MIT
- Density at separatrix ~ 1e20 m⁻³
- Power into lower-half domain 1-5 MW

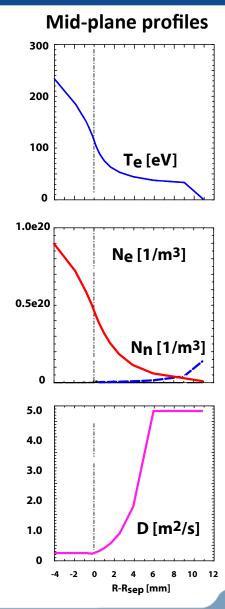
A C-Mod like case with ADX parameters is used for comparison



- Two configurations XPTD and SVPD
- Same underlying magnetic geometry, physics model, boundary conditions, etc.
- In SVPD the legs cut short to roughly match C-Mod vertical plate configuration

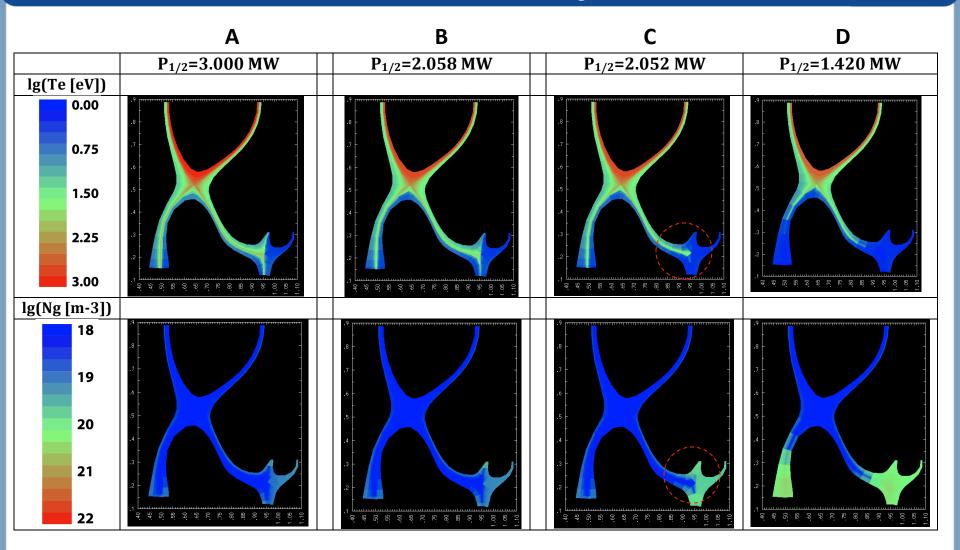
Radial transport parameters are set to match projected ADX upstream SOL characteristics

- Using fully recycling wall B.C. on all material surfaces
- Using radially growing diffusing coefficient to match the expected density profile width ~5 mm
- Spatially constant $\chi_{e,i}$ is sufficient to achieve ~3 mm width of mid-plane $T_{e,i}$
- Mid-plane profile projections are based on C-Mod data*

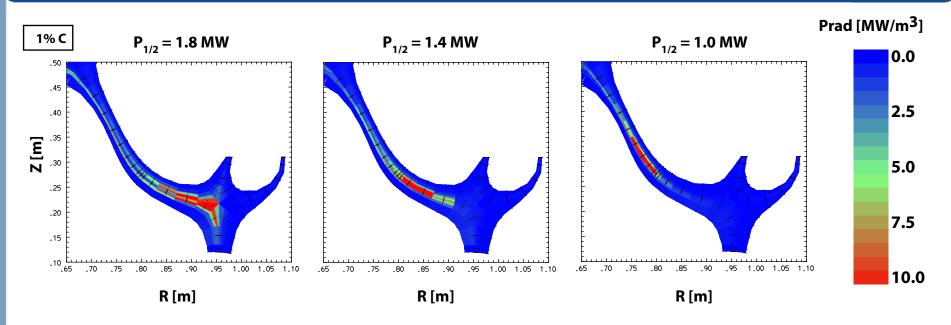


^{*}LaBombard et al., Nucl. Fusion 55, 053020, 2015.

As the input power $P_{1/2}$ is reduced, the divertor transitions to fully detached state

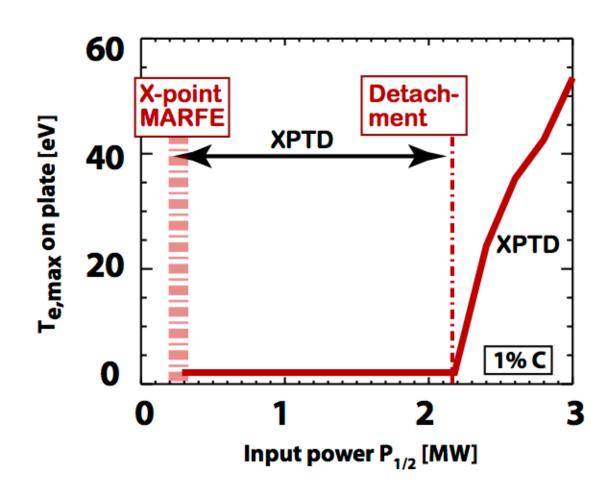


As input power P_{1/2} is reduced, radiation front remains stable but shifts upstream

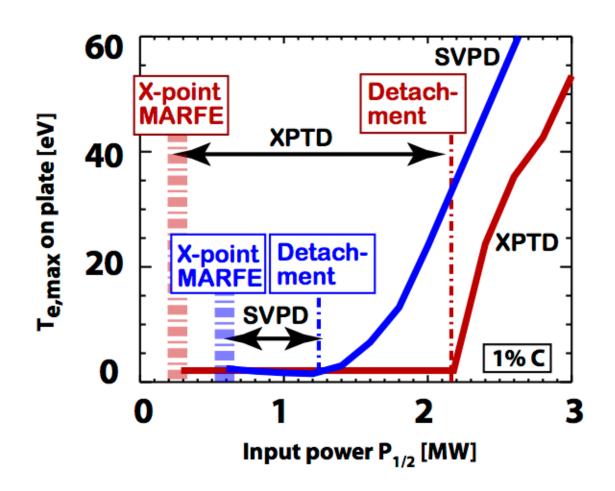


- For higher input power the radiation front moves to larger R to increase the radiating volume*
- P_{1/2} ≈ 2 MW onset of detachment
- As $P_{1/2}$ is reduced further either (i) the radiation front reaches the primary X-point, or (ii) no steady-state solutions can be found => X-point MARFE

Reducing input power P_{1/2} eventually leads to X-point MARFE



Reducing input power P_{1/2} eventually leads to X-point MARFE



Qualitatively similar results also obtained with 1% Ne impurity



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Summary, conclusions, plans

- Capability to generate grids with a secondary X-point is developed
- Capability to model configurations with a secondary X-point is developed in UEDGE
- Near-SNF configurations in DIII-D have been analyzed with UEDGE; points to strongly enhanced transport near the null (churning mode?)
- X-point Target Divertor (XPTD) configuration is studied with UEDGE for parameters matching the design of ADX tokamak
 - Steady state detachment found for XPTD, for a range of parameters
 - Easier to achieve detachment than for short leg divertor
 - Detachment front stays far away from the main X-point
 - Stable fully detached regimes for tokamak divertor?